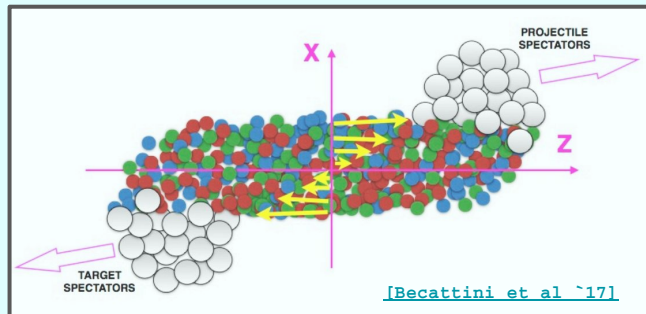


# Hydrodynamics of Globally Rotating Fluids Via Holography

Markus A. Garbiso



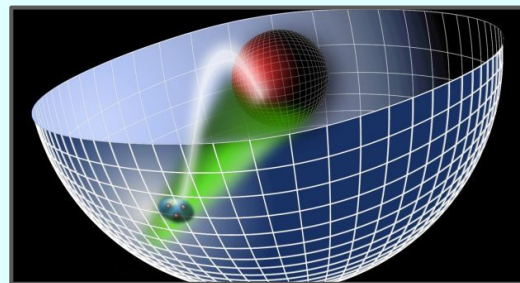
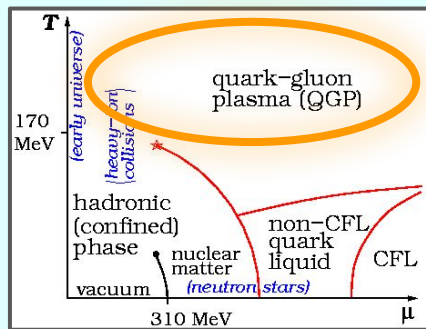


# Outline

- [Understanding the Field](#)
- [Holography and Why?](#)
- [Current Research - Globally Rotating Fluid](#)
- [Current Research - QNMs](#)
- [Current Research - Hydrodynamics](#)
- [Current Results -  \$\eta\_{++}\$](#)
- [Current Open Problem and Discussion](#)

# Understanding the Field

- Holography (AdS(Anti-de Sitter)/CFT Duality [\[Maldacena '99\]](#))
- QGP
  - RHIC [\[Braun-Munzinger et al. '01\]](#)
  - Vorticity Measured [\[The STAR Collaboration '17\]](#)
- Strongly Coupled Systems with **Broken Symmetries**
  - [Non-relativistic Systems \(Hořava Gravity\)](#) [\[Garbiso et al '19\]](#)
  - Systems with strong magnetic fields [\[Cartwright et al '19\]](#) [\[Ammon et al '17\]](#)
  - Rotating fluids from holography? [\[Hawking '99\]](#) [\[Reall et al '99\]](#)



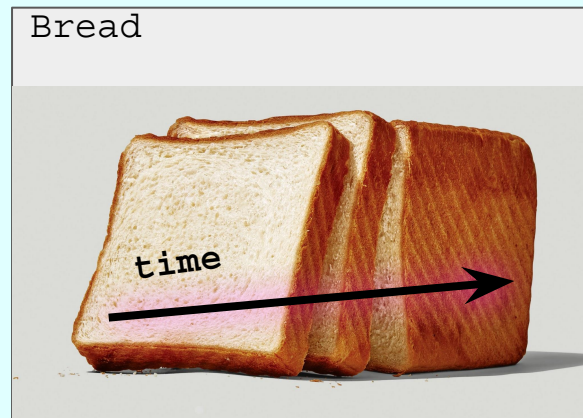


# Holography and Why?

- Concrete realization of the holographic principle: AdS/CFT
  - Strongest Statement: Type IIB Superstring theory  $AdS_5 \times S^5$  **"Gravity"** /  $\mathcal{N}=4$  SYM w/  $SU(N)$  **"CFT"** [\[Maldacena '99\]](#)
  - Weakest Statement (large N limit): *Supergravity*  $AdS_5 \times S^5$  **"Gravity"** / *Strongly Coupled QFT* **"CFT"**
  - Gravity/Fluid (large N limit) [\[Bhattacharyya et al '08\]](#) [\[Bhattacharyya '09\]](#)
  - Gravity/QGP [\[Busza et al '18\]](#)
- "Little Bangs" at RHIC
  - Understanding the fluid highly vortical **strongly coupled** fluids (QGP) demands explanation. [\[Busza et al '18\]](#)
- Mysteries of Higher Dimensional Gravity (Final State of Kerr AdS)

# Past Research - Hořava

- Using Holography: Non-Relativistic Gravity Theory
- Non-Relativistic Gravity Theory
  - Hořava Gravity [\[Hořava '09\]](#)  $t \rightarrow \tilde{t}(t), x^i \rightarrow \tilde{x}^i(t, x^i)$
  - Einstein-Aether Gravity [\[Jacobson et al '01\]](#)
  - Equivalent Theories [\[Bhattacharyya '13\]](#)
  - Simultaneity and Causality
- **(3+1)D** black brane: Dual to a **(2+1)D** sheet of non-relativistic material



# Past Research - Einstein- Aether Theory

$$S_{\text{aether}} = \frac{1}{4\pi G_{ae}} \int d^4x \sqrt{-g} \left( R - 2\Lambda + c_4 (u^\mu \nabla_\mu u^\nu) (u^\sigma \nabla_\sigma u_\nu) \right. \\ \left. - c_3 (\nabla_\mu u^\nu) (\nabla_\nu u^\mu) - c_2 (\nabla_\mu u^\mu)^2 - c_1 (\nabla^\mu u^\nu) (\nabla_\mu u_\nu) \right)$$

[\[Janiszewski '15\]](#)  
[\[Bhattacharyya '13\]](#)

aether vector field and khronon scalar field

$$u_\mu = \frac{\partial_\mu \phi}{\sqrt{-(\partial_\nu \phi)(\partial^\nu \phi)}}$$

coupling constants

$$\alpha = \frac{c_4}{1-c_3} = 0 \quad 1 + \lambda = \frac{1+c_2}{1-c_3} \quad \frac{G_H}{G_{ae}} = 1 + \beta = \frac{1}{1-c_3}$$

# Previous Research RESULTS

- Is this black brane a linearly stable spacetime?

- Yes, more [next slide](#)

- How do you perturb in [this theory](#)?

- $\longrightarrow$

$$g_{\mu\nu}^p = g_{\mu\nu} + \epsilon h_{\mu\nu}(x^\sigma)$$

- Are there scale invariances?

- YES wrt  $\beta$  and  $r_h$

$$\phi^p = \phi + \epsilon \chi(x^\sigma),$$

- Any connections to the relativistic version?

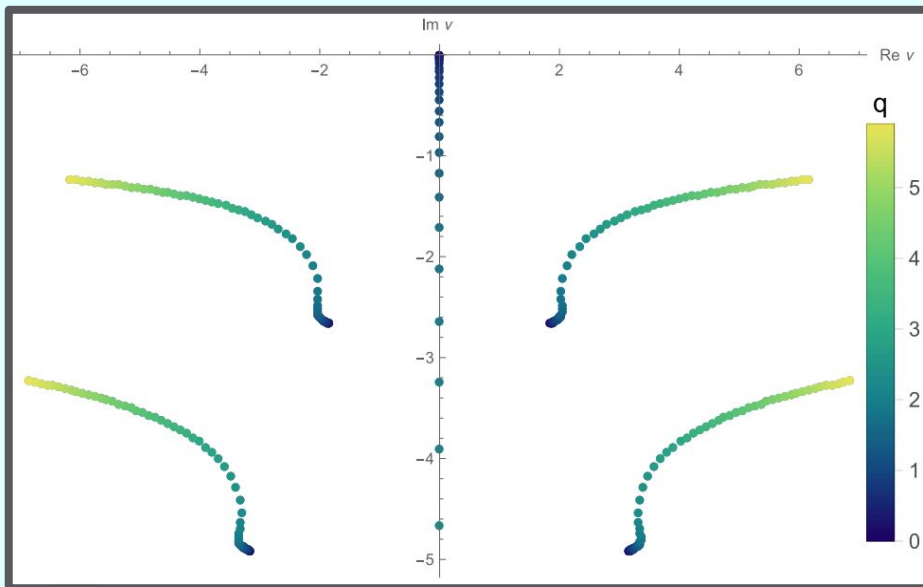
- Yes  $\omega_{\text{Horava}}^{\text{axial}} = \sqrt{1+\beta} \omega_{\text{Einstein}}^{\text{axial}}$

- and  $\omega_{\text{Horava}}^{\text{polar}} = \sqrt{1+\beta} \omega_{\text{Einstein}}^{\text{polar}}$  **MOSTLY**

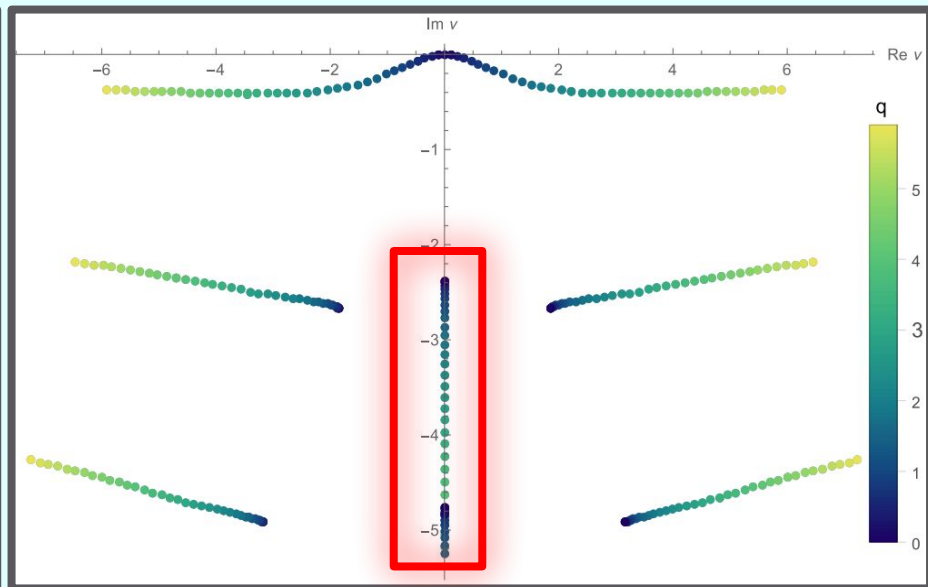
$$\nu = r_h \omega / (2^{1/3} \sqrt{1+\beta})$$

# Previous Research RESULTS

Axial QNMs



Polar QNMs



  = "Aether Modes"

$q$  = momentum, starting at 0

$$\nu = r_h \omega / (2^{1/3} \sqrt{1 + \beta})$$





# Current Research - Globally Rotating Fluid

- RHIC & QGP
  - “Non-central collisions have angular momenta of the order of  $1,000\hbar$ , and the resulting fluid may have a strong vortical structure 2–4 that must be understood to describe the fluid properly.” [\[The STAR Collaboration ‘17\]](#)
  - Modeling QGP plasma with strong vortical effects where effects are purely rotational
- Gauge/gravity and fluid/gravity
  - A new perspective on rotating strongly coupled fluids
  - How is bulk spacetime rotation encoded on the boundary fluid/field?
  - More QNMs, linear stability analysis?
- Final “AdS-Kerr State”?



## Current Research - Theory

$$S = -\frac{1}{16\pi} \int d^5x \sqrt{-g} \left( R + \frac{12}{L^2} \right)$$

5D MP AdS

$$ds^2 = - \left( 1 + \frac{r^2}{L^2} \right) dt^2 + \frac{dr^2}{G(r)} + \frac{r^2}{4} ((\sigma^1)^2 + (\sigma^2)^2 + (\sigma^3)^2) + \frac{2\mu}{r^2} \left( dt + \frac{a}{2} \sigma^3 \right)^2$$

$$G(r) = 1 + \frac{r^2}{L^2} - \frac{2\mu(1 - a^2/L^2)}{r^2} + \frac{2\mu a^2}{r^4}$$

$$\sigma^1 = -\sin \psi d\theta + \cos \psi \sin \theta d\phi$$

$$\sigma^2 = \cos \psi d\theta + \sin \psi \sin \theta d\phi$$

$$\sigma^3 = d\psi + \cos \theta d\phi$$

# Current Research - Theory (TBD) [\[Murata '09\]](#)

$$ds^2 = - (1 + r^2) dt^2 + \frac{dr^2}{G(r)} + \frac{r^2}{4} (4\sigma^+ \sigma^- + (\sigma^3)^2) + \frac{2\mu}{r^2} \left( dt + \frac{a}{2} \sigma^3 \right)^2.$$

$$\sigma^\pm = \frac{1}{2} (\sigma^1 \mp i\sigma^2)$$

$$e_\pm = \frac{1}{2} (e_1 \pm ie_2)$$

$$\sigma^a e_b = \delta_b^a$$

# Current Research - Theory (TBD) [\[Murata '09\]](#)

$$\begin{aligned}
 \xi_x &= \cos \phi \partial_\theta + \frac{\sin \phi}{\sin \theta} \partial_\psi - \cot \theta \sin \phi \partial_\phi \\
 \xi_y &= -\sin \phi \partial_\theta + \frac{\cos \phi}{\sin \theta} \partial_\psi - \cot \theta \cos \phi \partial_\phi \\
 \xi_z &= \partial_\phi
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{L}_{\xi_\alpha}(\sigma^a) &= 0 \\
 R_t \times SU(2) \times U(1) &\cong R_t \times U(2) \\
 \mathcal{L}_{\partial_t}(ds^2) &= 0 \\
 \mathcal{L}_{e_3}((\sigma^1)^2 + (\sigma^2)^2) &= 0 \\
 \mathcal{L}_{e_3}(\sigma^3) &= 0 \\
 e_3 &= \partial_\psi \\
 \mathcal{L}_{e_3}(\sigma^\pm) &= \pm \sigma^\pm
 \end{aligned}$$



# Current Research - Theory (TBD) [\[Murata `09\]](#)

$$L_\alpha = i\xi_\alpha, \quad W_a = ie_a, \quad \sum_{\alpha=x}^z L_\alpha L_\alpha = L^2 = W^2 = \sum_{a=1}^3 W_a W_a$$

$$[L_\alpha, L_\beta] = i\epsilon_{\alpha\beta\gamma} L_\gamma, \quad [W_a, W_b] = -i\epsilon_{abc} W_c, \quad [W_a, L_\alpha] = 0$$

# Current Research - Linear Analysis

Irreducible Representation of SU(2): **Wigner-D function**

$$L^2 D_{\mathcal{KM}}^{\mathcal{J}} = \mathcal{J}(\mathcal{J} + 1) D_{\mathcal{KM}}^{\mathcal{J}}, \quad L_z D_{\mathcal{KM}}^{\mathcal{J}} = \mathcal{M} D_{\mathcal{KM}}^{\mathcal{J}}, \quad W_3 D_{\mathcal{KM}}^{\mathcal{J}} = \mathcal{K} D_{\mathcal{KM}}^{\mathcal{J}}.$$

$$\mathcal{J} = 0, 1/2, 1, \dots \quad |\mathcal{K}| \leq \mathcal{J} \quad |\mathcal{M}| = 0 \leq \mathcal{J}$$

$$g_{\mu\nu}^p dx^\mu dx^\nu = (g_{\mu\nu} + \epsilon h_{\mu\nu} + O(\epsilon^2)) dx^\mu dx^\nu$$

$$\dot{R}_{\mu\nu} = \frac{2\Lambda}{D-2} h_{\mu\nu} \quad \dot{R}_{\mu\nu} = -\frac{1}{2} \nabla_\mu \nabla_\nu h - \frac{1}{2} \nabla^\lambda \nabla_\lambda h_{\mu\nu} + \nabla^\lambda \nabla_{(\mu} h_{\nu)\lambda}$$

[\[Wald `84\]](#)

$$\text{ex. } \phi = \sum_{\mathcal{K}} \phi^{\mathcal{K}}(r, t) D_{\mathcal{KM}}^{\mathcal{J}}(\theta, \phi, \psi)$$

**perturbations of different  $((\mathcal{J}, \mathcal{M}), \mathcal{K})$  decouple** [\[Murata et al `08\]](#)

# Current Research - h++ Modes (TBD)

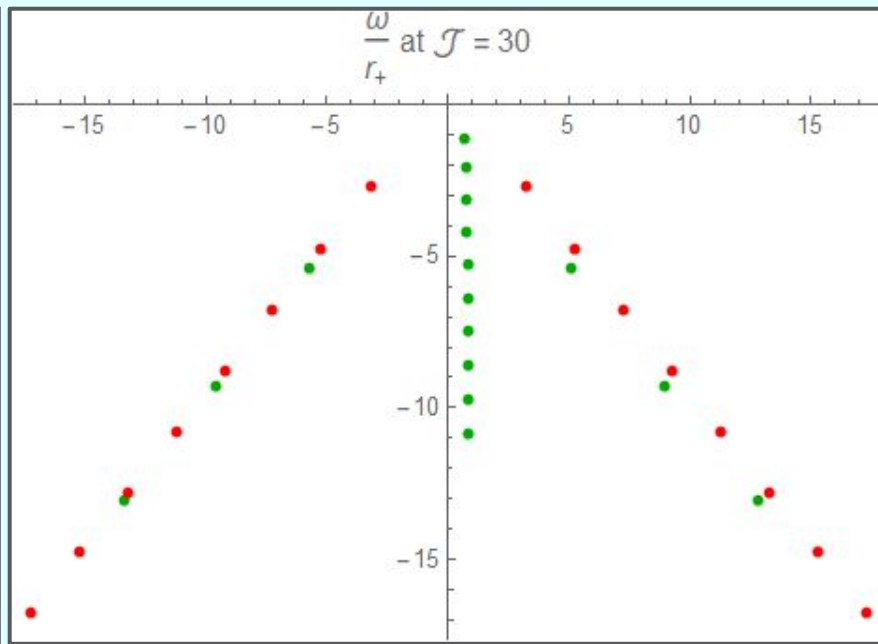
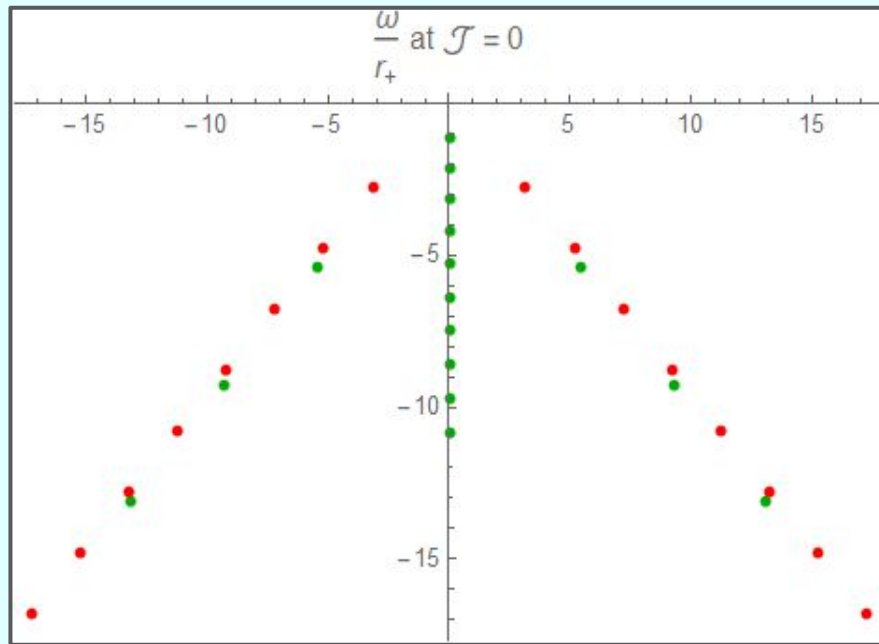


# Current Results - QNMs - “h++” - Real vs Im

At large temperatures ie.  $r_+/L \sim 100$

● 90% Extremality

● No Rotation



**Note new modes close to imaginary axis from larger  $a$ .**

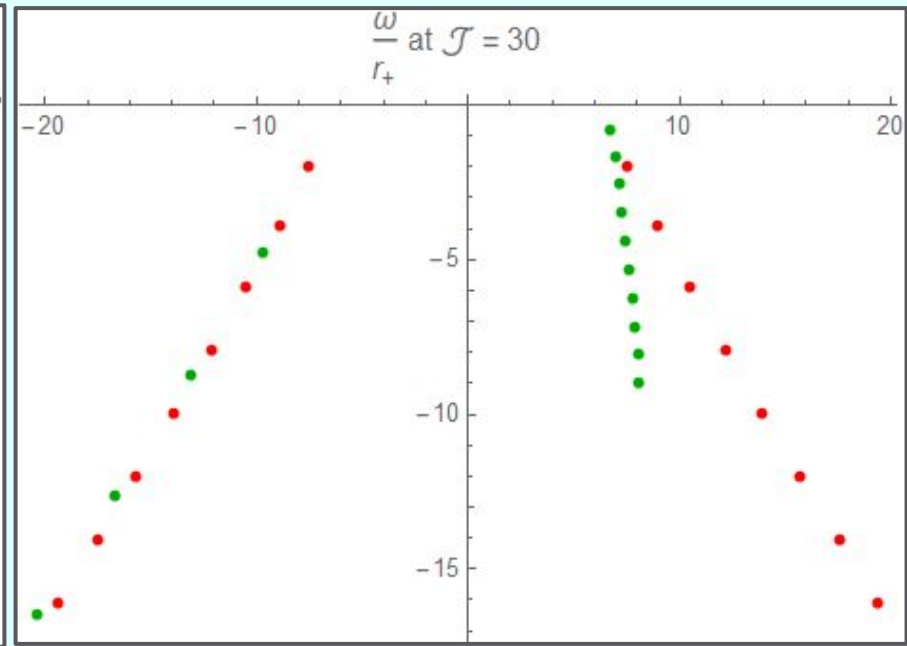
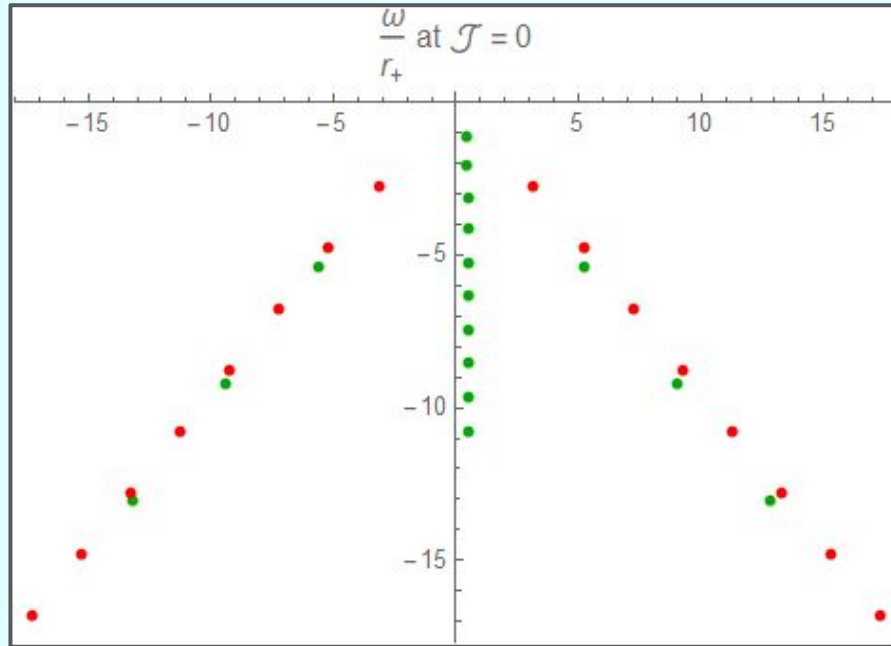


# Current Results - QNMs - “h++” - Real vs Im

At small temperatures ie.  $r_+/L \sim 10$

● 90% Extremality

● No Rotation

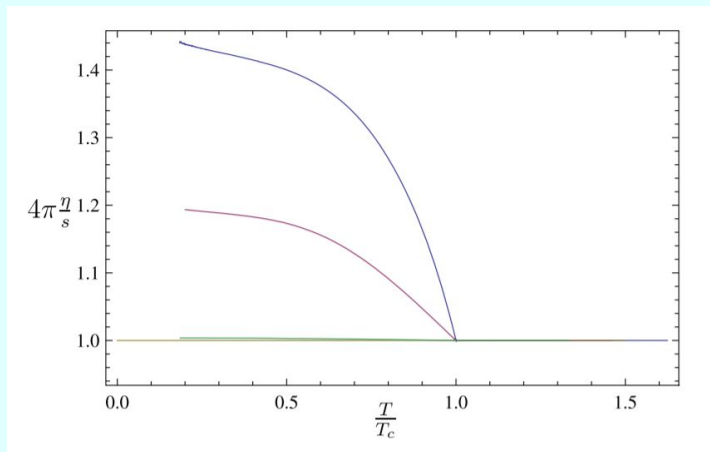


Note that modes are not symmetric across Imaginary axis.  
Note higher  $J$  modes are more “unstable”.

# Current Research - Hydrodynamics [\[Cardoso et al. `14\]](#)

Rotation seemingly distinguishes a direction so does the **shear viscosity**,  $\mu$ , change? Shear viscosity doesn't respect the "1/4 $\pi$ " if anisotropy is induced.

[\[Erdmenger et al. `11\]](#) [\[Critelli et al `14\]](#)



$$S = \frac{1}{2\kappa_5^2} \int d^5x \left( (R - 2\Lambda) - \frac{\alpha^2}{2} \text{Tr}(F^2) \right)$$

[\[Erdmenger et al. `11\]](#)



# Current Results - Needed Calculations

Now, let's find the temperature of our black hole.

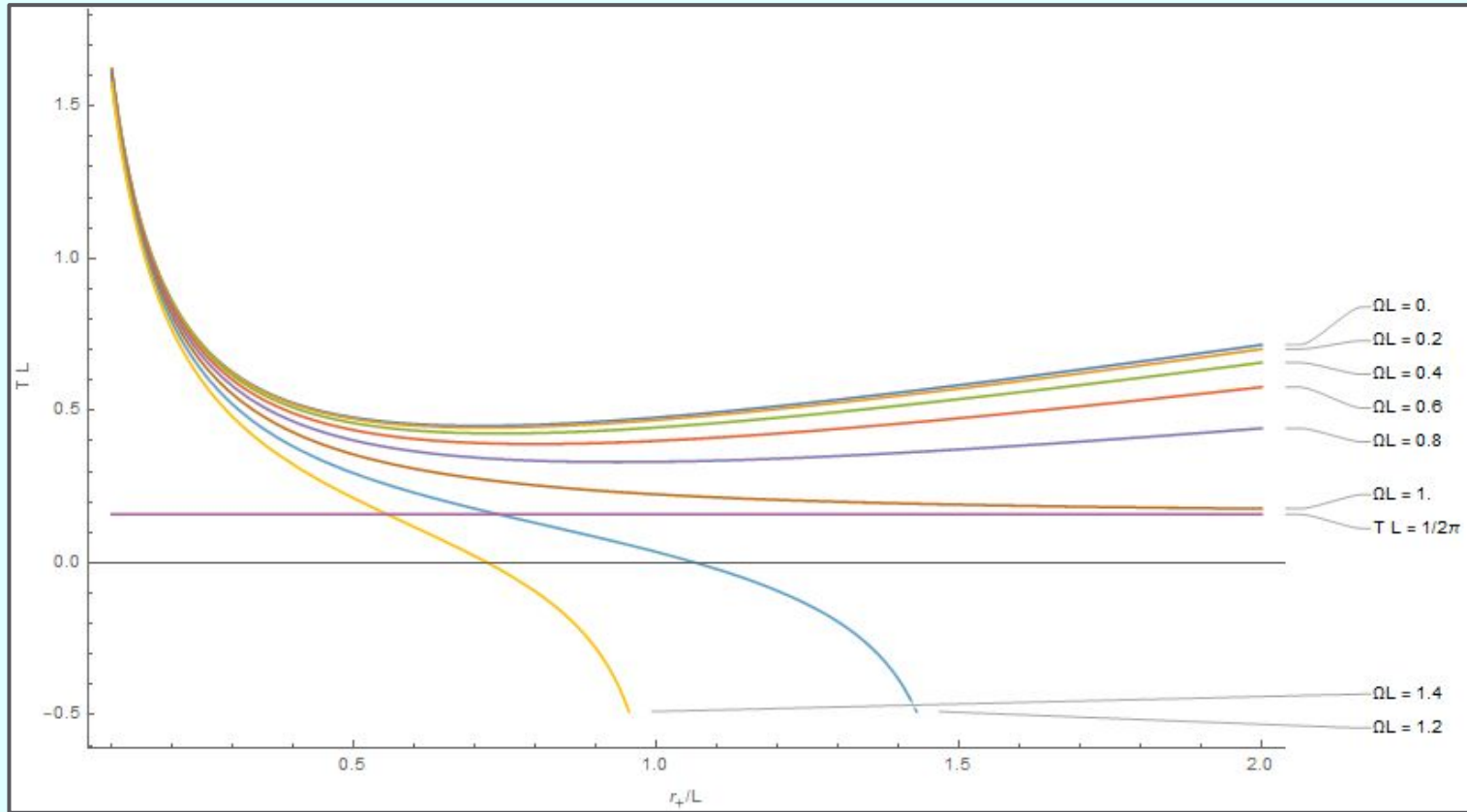
$$T = \kappa/2\pi = \sqrt{\frac{r_+^2 + 1}{r_+^2 (1 - \Omega^2) + 1}} \frac{2r_+^2 (1 - \Omega^2) + 1}{2\pi r_+}$$

Hydrodynamics requires large temperature limits and therefore implies " $\Omega L < 1$ ".

Furthermore, superradiant instabilities no longer apply [\[Cardoso et al. '14\]](#) [\[Murata '09\]](#), because extremality limit and faster than light limit coincide for .



# Hydrodynamics - Large Temperature



# Current Results - $\eta_{++}$ (in $L=1$ units)

Kubo Formula for shear viscosity

From papers [\[Son et al '02\]](#), we conjectured the expression to the right as a possible expression for  $\eta$ .

$$\eta = - \lim_{\bar{\omega} \rightarrow 0^+} \frac{\text{Im} \langle T_{++} T_{++} \rangle}{\bar{\omega}/T} = - \lim_{\bar{\omega} \rightarrow 0^+} \frac{\text{Im} h_{++}^{(4)} / h_{++}^{(0)}}{\bar{\omega}/T}$$

$$\sim - \frac{T \text{ vev}}{\bar{\omega} \text{ source}} \Big|_{\bar{\omega}=\text{small}}$$

# Current Results - $\eta_{++}$ (recap)

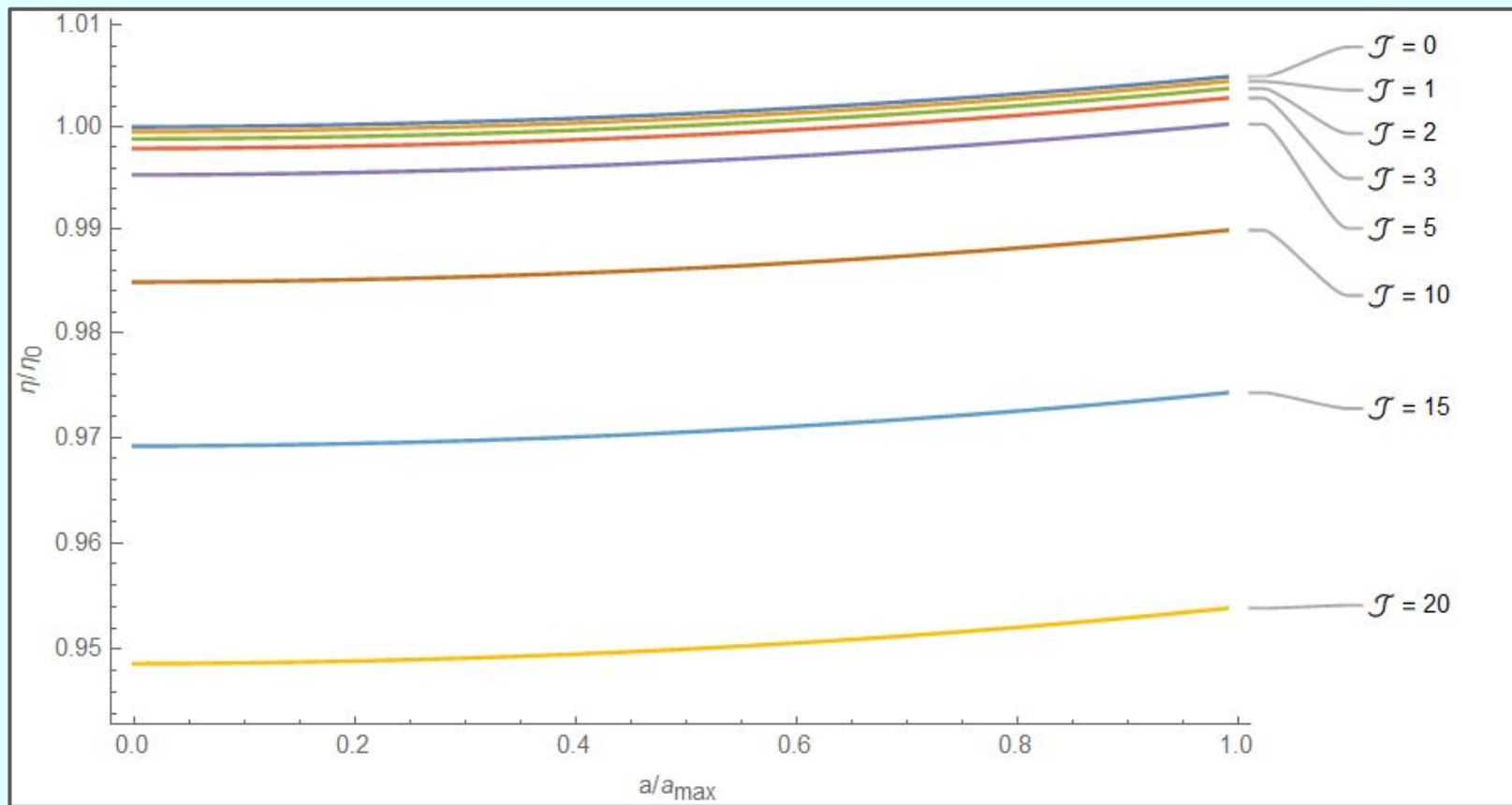
$$T = \kappa/2\pi = \sqrt{\frac{r_+^2 + 1}{r_+^2 (1 - \Omega^2) + 1}} \frac{2r_+^2 (1 - \Omega^2) + 1}{2\pi r_+}$$

$$G(r_+) = 0 \qquad \chi^2|_{r=r_+} = 0$$

$$\chi = \partial_t - 2\Omega \partial_\psi$$

$$\eta \sim - \frac{T_{\text{vev}}}{\bar{\omega}_{\text{source}}} \Big|_{\bar{\omega}=\text{small}}$$

# Current Results - $\eta_{++}$ @ ( $r_+ = 100$ )



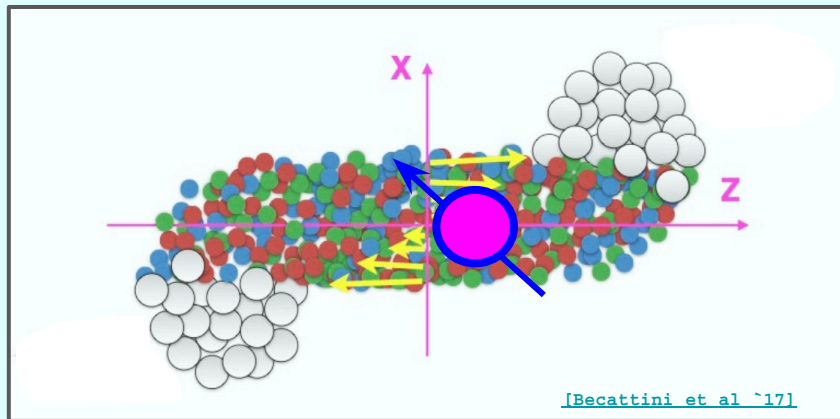
# Current Open Problems

How does one model **hyperon**-like fields [\[Florkowski '19\]](#)

QNMs of probe fields

- positive integer spins?
- orbital spin & spin coupling?
- spinors

Finding  $\eta$  legitimately.





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I would also like to thank **Ochanomizu University, The University of Tokyo, Nihon University,** and **Kyoto University** for allowing me to visit and present my research at their campuses.





**Any  
Questions?**



# Useful Refs

<https://arxiv.org/pdf/1301.2826.pdf> Flow and Viscosity in Relativistic Heavy Ion Collisions  
<https://arxiv.org/pdf/1905.11309.pdf> Quasinormal modes for quarkonium in a plasma with magnetic fields  
<https://arxiv.org/pdf/1308.2672.pdf> Kerr-AdS analogue of triple point and solid/liquid/gas phase transition  
<https://arxiv.org/pdf/1401.2586.pdf> Thermodynamics of rotating black holes and black rings: phase transitions and thermodynamic volume  
<https://arxiv.org/pdf/1510.04713.pdf> Maxwell perturbations on asymptotically anti-de Sitter spacetimes: Generic boundary conditions and a new branch of quasinormal modes  
<https://arxiv.org/abs/1312.5323> Holographic thermalization, quasinormal modes and superradiance in Kerr-AdS  
<https://arxiv.org/pdf/1505.04793.pdf> Black holes with a single Killing vector field: black resonators  
<https://arxiv.org/pdf/0803.1371.pdf> Stability of Five-dimensional Myers-Perry Black Holes with Equal Angular Momenta  
<https://arxiv.org/pdf/0901.2574.pdf> Warped AdS5 Black Holes and Dual CFTs  
<https://arxiv.org/pdf/1302.1580.pdf> Boundary Conditions for Kerr-AdS Perturbations  
<https://arxiv.org/pdf/1802.04801.pdf> Heavy Ion Collisions: The Big Picture, and the Big Questions  
<https://arxiv.org/pdf/0904.2154.pdf> Gravitational stability of simply rotating Myers-Perry black holes: tensorial perturbations  
<https://www.sciencedirect.com/science/article/pii/S0370269311003959?via%3Dihub> Non-universal shear viscosity from Einstein gravity  
<https://arxiv.org/pdf/0803.1371.pdf> Stability of Five-dimensional Myers-Perry Black Holes with Equal Angular Momenta