

Since the topic of index theorems contains many pieces of mathematical vocabulary here are some practice problems to help you learn all of the words I used in seminar.

1) What is the polynomial symbol associated with the following operators,

$$1. \frac{d^2 f}{dx^2} = 0$$

$$2. \frac{d^2 f}{dx^2} - \frac{d^2 f}{dy^2} = 0$$

$$3. \frac{d^2 f}{dx^2} + \frac{d^2 f}{dy^2} + 8 \frac{d^2 f}{dy dx} = 0$$

Are these operators elliptic? Additionally show that the following differential operator defined via the symbol

$$\sigma(\nabla, \xi) = A_{11}\xi_1^2 + A_{21}\xi_1\xi_2 + A_{22}\xi_2^2 \tag{1}$$

is elliptic if and only if the symbol σ is an ellipse in ξ space. I'll remind you that to show an if and only if statement you must show both directions. That is suppose x is y if and only if y is z means suppose x is y show that y is z , and suppose y is z show x is y .

2) The special linear group is defined

$$SL(2, \mathbb{C}) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid a, b, c, d \in \mathbb{C} \text{ and } ab - bc = 1 \right\} \tag{2}$$

Notice that this is not the special unitary group $SU(2)$ in fact $SU(2) \subset SL(2, \mathbb{C})$. Additionally let $SO(1, 3)^\uparrow$ be the special orthogonal group with $\Lambda_0^0 > 0$ i.e. the proper orthochronous Lorentz group. It turns out that the special linear group is a double cover of the proper orthochronous Lorentz group. This is just what we did in group theory this last week. So you know the drill, given that the homomorphism is $\Lambda : SL(2, \mathbb{C}) \rightarrow SO(1, 3)^\uparrow$ is given by $\Lambda(g)X = gXg^\dagger$ show that the action $\Lambda(g)X$ with,

$$g = \begin{pmatrix} \cosh \alpha/2 + \sinh \alpha/2 & 0 \\ 0 & \cosh \alpha/2 - \sinh \alpha/2 \end{pmatrix} \tag{3}$$

is the same as a boost around the z -axis. The velocity will be $\nu = \tanh \alpha$. Hints: write $\Lambda(g)X = x' = x'^\mu \sigma_\mu$ with $\sigma_\mu = (I, \vec{\sigma})$ and call x'^μ the transformed vector. Then write $X(x) = x^\mu \sigma_\mu$ and compute gXg^\dagger for the given element. Then $x'^\mu \sigma_\mu = gX(x)g^\dagger$.

3) In preparation for next week consider the following action,

$$\int_0^\beta dt \left(\frac{1}{2} g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu + \frac{1}{2} g_{\mu\nu} \psi^\mu \frac{D\psi^\nu}{Dt} \right) \tag{4}$$

Here $\frac{D\psi^\nu}{Dt} = \frac{d\psi^\nu}{dt} + \dot{x}^\lambda \Gamma_{\lambda\kappa}^\nu \psi^\kappa$ is the covariant derivative along the curve $x(t)$. Calculate the equations of motion for ψ and for x . Assuming that ψ is a grassman number (i.e. $\theta\phi = -\phi\theta$ for grassman numbers ϕ and θ) and x is a classical (commuting) number.

4) Speaking of grassman numbers have you ever calculated the integral of a grassman number before? If not this problem will fix you right up. These two formulas will be good to know.

$$\int d\theta = 0, \quad \int d\theta\theta = 1 \quad (5)$$

That's right. To calculate the value you take the derivative of the integrand. With this in mind the fermionic harmonic oscillator is described by grassman numbers. The Hilbert space is spanned by $\mathcal{H} = \text{span}\{|0\rangle, |1\rangle\}$. Knowing that $\theta^2 = 0$ and $|\theta\rangle = |0\rangle + |1\rangle\theta$, show that the completeness relation can be written,

$$\int d\theta^* d\theta |\theta\rangle\langle\theta| e^{-\theta^*\theta} = I \quad (6)$$